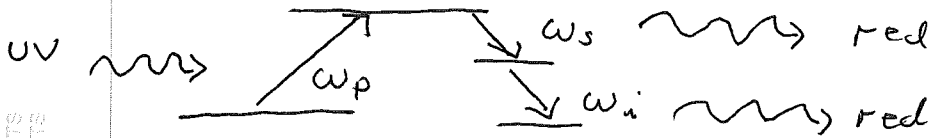


7.2 The SQUEEZE GENERATION

SPDC = spontaneous parametric down conversion $\Rightarrow \square \Rightarrow$

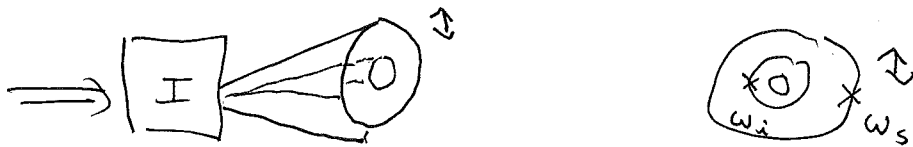
OPA = optical parametric amplifier $\rightarrow \square \rightarrow$

OPO = " " oscillator $\rightarrow (\square) \rightarrow$



To conserve energy: $\hbar\omega_p = \hbar\omega_s + \hbar\omega_i$
 To conserve momentum: $\hbar\vec{k}_p = \hbar\vec{k}_s + \hbar\vec{k}_i$

Type I Polarization all the same



Type II Polarization ⊥



Let $\omega_s = \omega_i = \omega = \omega_p/2$

$$\hat{H} = \underbrace{\hbar\omega\hat{a}^\dagger\hat{a}}_{\text{signal}} + \underbrace{\hbar\omega_p\hat{b}^\dagger\hat{b}}_{\text{pump}} + i\hbar\chi^{(2)} \left[\underbrace{\hat{a}^2\hat{b}^\dagger}_{\substack{\nearrow \omega_p \\ \searrow \omega \\ \searrow \omega}} - \underbrace{\hat{a}^{\dagger 2}\hat{b}}_{\substack{\nearrow \omega_p \\ \searrow \omega \\ \nearrow \omega}} \right]$$

$\chi^{(2)}$ is second order nonlinearity

$$\vec{P} = \underbrace{(\chi^{(1)} + \chi^{(1)})}_{n} \vec{E} + (\chi^{(2)} \vec{E}) \vec{E} + (\chi^{(3)} E^2) \vec{E} \dots$$

Expansion of polarization \vec{P} of a nonlinear dielectric

SHO $\vec{F}(x) = \underbrace{kx}_{\text{linear}} + \underbrace{\alpha x^2}_{x^{(2)}} + \underbrace{\beta x^3}_{x^{(3)}} + \dots$
 anharmonic oscillator

Parametric nondepleted pump approximation

$\hat{b} \rightarrow \beta e^{-i\omega_p t} \Rightarrow \hat{b}^\dagger \rightarrow \beta^* e^{i\omega_p t}$

Pump is ∞ reservoir of $\hbar\omega_p$ photons
 (Typical: Watts \rightarrow kilowatts)

$\Rightarrow \hat{H} = \hbar\omega_p \hat{n} + \hbar\omega_p |\beta|^2 + i\hbar\alpha^{(2)} [\hat{a}^2 \beta^* e^{i\omega_p t} - \hat{a}^{\dagger 2} \beta e^{-i\omega_p t}]$
 constant does not contribute

Recall Schrödinger \rightarrow Heisenberg
 $\hat{a} \rightarrow \hat{a}(t) = \hat{a} e^{-i\omega t}$

$\Rightarrow \hat{H}_{\text{Interaction}} = i\hbar\alpha^{(2)} [\hat{a}^2 \beta^* e^{i(\omega_p - 2\omega)t} - \hat{a}^{\dagger 2} \beta e^{-i(\omega_p - 2\omega)t}]$

But $\omega_p = \omega_s + \omega_i = 2\omega$ Eng. Cons.

So $\omega_p - 2\omega = 0$ Let $\zeta \equiv \alpha^{(2)} \beta$

$\Rightarrow \hat{H}_I = i\hbar [\zeta^* \hat{a}^2 - \zeta \hat{a}^{\dagger 2}]$

Since \hat{H}_I is time independent, time evolution operator is (see e.g. Merzbacher)

$\hat{U}_I(t, 0) \equiv e^{-i\int \hat{H}_I dt}$

Hence

$$\begin{aligned}
 U_I(t) &= e^{-i\gamma t [i\hbar(\hat{a}^* \hat{a}^2 - \hat{a} \hat{a}^{\dagger 2})]} \\
 &= \exp \left[\frac{2\hat{a}^* t}{2} \hat{a}^2 - \frac{2\hat{a} t}{2} \hat{a}^{\dagger 2} \right] \\
 &= \exp \left[\frac{\zeta^*(t)}{2} \hat{a}^2 - \frac{\zeta(t)}{2} \hat{a}^{\dagger 2} \right] \\
 &= \hat{S}(\zeta(t)) \quad \therefore
 \end{aligned}$$

Hence $\boxed{\zeta = 2\hat{a} t = 2\beta \chi^{(2)} t}$

is the squeezed param

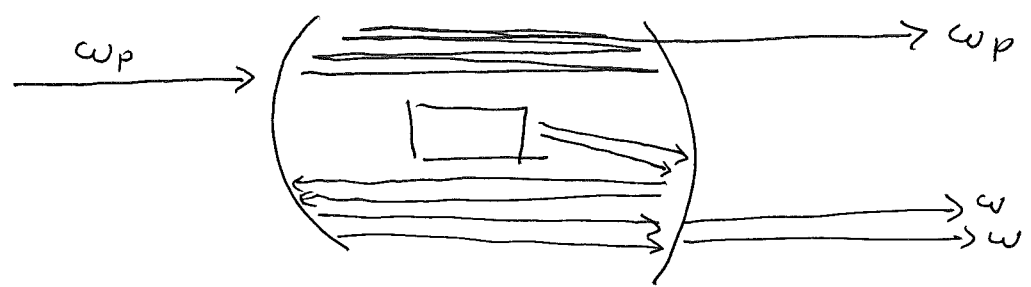
To maximize $\zeta \gg 1$

$\Rightarrow \beta \gg 1$ STRONG PUMP AMPLITUDE

OR $t \gg 1$ LONG INTERACTION TIME
(USE LONG XTAL OR PUT XTAL IN FABRY PEROT)

OR $\chi^{(2)} \gg 1$ CHEMISTRY! CHINESE MAKE XTALS WITH LARGE $\chi^{(2)}$!

OPO

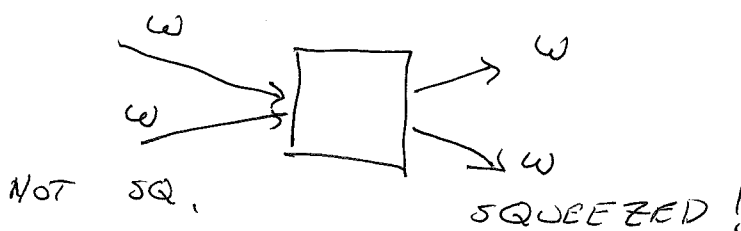
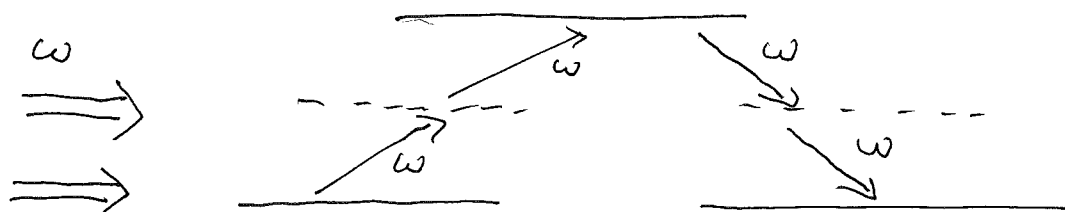


OPO make mirrors at ω_p or ω or both!

22-102 50 SHEETS
 22-102 100 SHEETS
 22-102 200 SHEETS

Four Wave Mixing

We can also use χ_3 for squeezing
 Less efficient ~~but~~ ($\chi_3 \ll \chi_2$)
 but advantage for quantum lithography that
 $(\omega_p, \omega_p) \rightarrow (\omega_p, \omega_p)$ no change in ω



$$\hat{H}_{\chi_3} = \underbrace{\hbar\omega a^\dagger a + \hbar\omega b^\dagger b}_{\text{same}} + i\hbar\chi_3 [a^{\dagger 2} b^{\dagger 2} - a^{+2} b^2]$$

Again assume strong classical parametric
 nondepleted pump

$$\Rightarrow \hat{b} = \beta e^{-i\omega t}$$

Interaction picture $\hat{a} \rightarrow \hat{a}(t) = \hat{a} e^{-i\omega t}$

$$\Rightarrow \hat{H}_I^{\chi_3} = i\hbar\chi_3 [\beta^{*2} \hat{a}^2 - \beta^2 \hat{a}^{\dagger 2}]$$

$$= \boxed{ i\hbar [\xi_3^* \hat{a}^2 - \xi_3 \hat{a}^{\dagger 2}] }$$

This is identical to previous result with

$$\xi_2 = \chi_2 \beta \rightarrow \xi_3 = \chi_3 \beta^2$$

22-141 50 SHEETS
22-142 100 SHEETS
22-144 200 SHEETS



Hence

$$e^{i4 \frac{\chi_3}{\epsilon} t} = \frac{1}{S_{\chi_3}(t)}$$

where $S_{\chi_3}(t) = 2\chi_3 \beta^2 t$ is FWMix

squeeze param.

Again want

$$\chi_3 \gg 1$$

$$\beta^2 \gg 1$$

$$t \gg 1$$

note $\chi_3 \ll \chi_2$ but $|\beta|^2 \gg |\beta|$

so with strong pump squeezing nearly as good. First squeezing done w/ χ_3 (silica)

χ_2

Crystal symmetry
Must be non axisymmetric

hard

special birefringent
XTALS

χ_3

any material
quartz
glass

easy