

R. F. O'CONNELL

21 Agosto 1969

Lettere al Nuovo Cimento

Serie I, Vol. 2, pag. 221-222

Simple Derivation of Schwinger's Quantization Relation between Electric and Magnetic Charges.

R. F. O'CONNELL

Department of Physics and Astronomy, Louisiana State University - Baton Rouge, La.

(ricevuto il 27 Giugno 1969)

DIRAC⁽¹⁾ has shown that the existence of magnetic charge (monopoles) g implies the quantization of electric charge e , i.e. $2ge = nhc$, where n is an integer. However, SCHWINGER⁽²⁾ has recently claimed that the correct relation is actually

$$(1) \quad ge = nhc.$$

Schwinger's derivation is quite complicated so, with the intention of obtaining a straightforward derivation of the quantization condition, EPINGER⁽³⁾ analysed the solution of Schrödinger's equation for a charge e in a uniform magnetic field H produced by an infinite plane sheet of total magnetic charge g . He obtained $ge = (n + \frac{1}{2})hc$. However, his analysis is incomplete because he neglected electron spin (i.e. he considered the Landau⁽⁴⁾ contribution to the quantization of electron orbits in a plane perpendicular to H , but he neglected the Pauli⁽⁵⁾ contribution).

Consider the motion of a relativistic electron in a magnetic field $H = H_z$ produced by an infinite plane sheet of total magnetic charge g in the xy -plane. The energy eigenvalues resulting from a solution of Dirac's equation are^(6,7) (henceforth we take $\hbar = c = m = 1$, where m is the mass of the electron)

$$(2) \quad E^2 = 1 + P_x^2 + 2enH, \quad n = 0, 1, 2, \dots$$

(1) P. A. M. DIRAC: *Proc. Roy. Soc., A* **133**, 60 (1931); *Phys. Rev.*, **74**, 817 (1948).

(2) J. SCHWINGER: *Phys. Rev.*, **144**, 1087 (1964).

(3) H. J. EPINGER: *Note on the quantum theory of elastic and magnetic charges* (preprint).

(4) L. LANDAU: *Zeits. Phys.*, **64**, 629 (1930).

(5) W. PAULI: *Zeits. Phys.*, **41**, 81 (1927).

(6) M. H. JOHNSON and H. A. LITTMANN: *Phys. Rev.*, **76**, 323 (1949); **77**, 702 (1950).

(7) A. A. SOKOLOV: *Introduction to quantum electrodynamics* (U.S. Atomic Energy Commission: AEC-tr-4322, 1960).

Thus, the momentum P_{\perp} in the xy -plane is quantized according to the relation $P_{\perp}^2 = 2enH$. Now, in contrast to the classical motion, we cannot say that the electron trajectory in the xy -plane is a circle of radius R . This arises from the fact that there is an uncertainty in locating the centre of the electron orbit ⁽⁵⁾ with the result that there is a quadratic fluctuation superimposed upon the circular motion ^(6,7). However, we can still talk about the radius R of the equilibrium circle about which the fluctuations take place ⁽⁸⁾. This radius corresponds to the position of maximum charge density ⁽⁹⁾ and is given by ^(8,9)

$$(3) \quad R = (2n/eH)^{1/2}.$$

It follows that $R = P_{\perp}/eH$, in agreement ^(8,9) with the classical result. Thus, the magnetic flux F enclosed within these quantized electron orbits is

$$(4) \quad F = \pi R^2 H = \pi 2n/e.$$

But $F = 2\pi g$. Therefore,

$$(5) \quad ge = n,$$

which is Schwinger's ⁽²⁾ quantization relation. In addition, it is interesting to note that the magnetic flux is quantized in units of $2\pi\hbar c/e$, the original unit derived by LONDON ⁽¹⁰⁾ in his discussion of quantized flux through a superconducting ring.

⁽⁵⁾ Ref. (7), eq. (28.23).

⁽⁶⁾ J. J. KLEIN: *Rev. Mod. Phys.*, **40**, 523 (1968).

⁽⁷⁾ P. LONDON: *Superfluids*, vol. 1 (New York, 1950), p. 152.